25.1. Model: Balmer's formula predicts a series of spectral lines in the hydrogen spectrum. **Solve:** Substituting into the formula for the Balmer series,

$$
\lambda = \frac{91.18 \text{ nm}}{\left(\frac{1}{2^2} - \frac{1}{n^2}\right)} \Rightarrow \lambda = \frac{91.18 \text{ nm}}{1/2} = 410.3 \text{ nm}
$$

where $n = 3, 4, 5, 6, \ldots$ and where we have used $n = 6$. Likewise for $n = 8$ and $n = 10, \lambda = 389.0$ nm and $\lambda = 379.9$ nm.

25.2. Model: Balmer's formula predicts a series of spectral lines in the hydrogen spectrum. **Solve:** Balmer's formula is

$$
\lambda = \frac{91.18 \text{ nm}}{\left(\frac{1}{2^2} - \frac{1}{n^2}\right)} \qquad n = 3, 4, 5, 6, ...
$$

As $n \to \infty$, $1/n^2 \to 0$. Thus, $\lambda_{n \to \infty} = 4(91.18 \text{ nm}) = 364.7 \text{ nm}$.

25.3. Model: Balmer's formula predicts a series of spectral lines in the hydrogen spectrum. **Solve:** Using Balmer's formula,

$$
\lambda = 389.0 \text{ nm} = \frac{91.18 \text{ nm}}{\left(\frac{1}{2^2} - \frac{1}{n^2}\right)} \Rightarrow \frac{1}{4} - \frac{1}{n^2} = 0.2344 \Rightarrow n = 8
$$

25.4. Model: The angles of incidence for which diffraction from parallel planes occurs satisfy the Bragg condition.

Solve: The Bragg condition is $2d\cos\theta_m = m\lambda$, where $m = 1, 2, 3, ...$ For first and second order diffraction,

$$
2d\cos\theta_1 = (1)\lambda \qquad 2d\cos\theta_2 = (2)\lambda
$$

Dividing these two equations,

$$
\frac{\cos \theta_2}{\cos \theta_1} = 2 \Longrightarrow \theta_2 = \cos^{-1}(2\cos \theta_1) = \cos^{-1}(2\cos 68^\circ) = 41^\circ
$$

25.5. Model: The angles of incidence for which diffraction from parallel planes occurs satisfy the Bragg condition.

Solve: The Bragg condition is $2d\cos\theta_m = m\lambda$. For $m = 1$ and for two different wavelengths,

$$
2d\cos\theta_1 = (1)\lambda_1 \qquad 2d\cos\theta_1' = (1)\lambda_1'
$$

Dividing these two equations,

$$
\frac{\cos \theta'_1}{\cos \theta'_1} = \frac{\lambda'_1}{\lambda_1} \Longrightarrow \frac{\cos \theta'_1}{\cos 54^\circ} = \frac{0.15 \text{ nm}}{0.20 \text{ nm}} \Longrightarrow \theta'_1 = \cos^{-1}(0.4408) = 64^\circ
$$

25.6. Model: The angles corresponding to the various orders of diffraction satisfy the Bragg condition. **Solve:** The Bragg condition for $m = 1$ and $m = 2$ gives

$$
2d\cos\theta_1 = (1)\lambda \qquad 2d\cos\theta_2 = (2)\lambda
$$

Dividing these two equations,

$$
\cos \theta_1 = \frac{\cos \theta_2}{2} = \frac{\cos 45^{\circ}}{2} \Rightarrow \theta_1 = \cos^{-1} \left(\frac{\cos 45^{\circ}}{2} \right) = 69.3^{\circ}
$$

25.7. Model: The angles corresponding to the various diffraction orders satisfy the Bragg condition. **Solve:** The Bragg condition is $2d\cos\theta_m = m\lambda$, where $m = 1, 2, 3, ...$ The maximum possible value of *m* is the number of possible diffraction orders. The maximum value of $cos\theta_m$ is 1. Thus,

$$
2d = m\lambda \Rightarrow m = \frac{2d}{\lambda} = \frac{2(0.180 \text{ nm})}{(0.085 \text{ nm})} = 4.2
$$

We can observe up to the fourth diffraction order.

25.8. Model: Use the photon model of light.

Solve: The energy of the photon is

$$
E_{\text{photon}} = hf = h\frac{c}{\lambda} = (6.63 \times 10^{-34} \text{ Js}) \left(\frac{3.0 \times 10^8 \text{ m/s}}{500 \times 10^{-9} \text{ m}} \right) = 3.98 \times 10^{-19} \text{ J}
$$

Assess: The energy of a single photon in the visible light region is extremely small.

25.9. Model: Use the photon model of light.

Solve: The energy of the single photon is

$$
E_{\text{photon}} = hf = h \left(\frac{c}{\lambda} \right) = \frac{(6.63 \times 10^{-34} \text{ Js})(3.0 \times 10^8 \text{ m/s})}{1.0 \times 10^{-6} \text{ m}} = 1.99 \times 10^{-19} \text{ J}
$$

\n
$$
\Rightarrow E_{\text{mol}} = N_{\text{A}} E_{\text{photon}} = (6.023 \times 10^{23})(1.99 \times 10^{-19} \text{ J}) = 1.2 \times 10^5 \text{ J}
$$

Assess: Although the energy of a single photon is very small, a mole of photons has a significant amount of energy.

25.10. Model: Use the photon model of light.

Solve: The energy of a photon with wavelength λ_1 is $E_1 = hf_1 = hc/\lambda_1$. Similarly, $E_2 = hc/\lambda_2$. Since E_2 is equal to $2E_1$,

$$
\frac{hc}{\lambda_2} = 2\frac{hc}{\lambda_1} \Rightarrow \lambda_2 = \frac{\lambda_1}{2} = \frac{600 \text{ nm}}{2} = 300 \text{ nm}
$$

Assess: A photon with $\lambda = 300$ nm has twice the energy of a photon with $\lambda = 600$ nm. This is an expected result, because energy is inversely proportional to the wavelength.

25.11. Model: Use the photon model of light.

Solve: The energy of the x-ray photon is

$$
E = hf = h\left(\frac{c}{\lambda}\right) = \left(6.63 \times 10^{-34} \text{ Js}\right) \left(\frac{3.0 \times 10^8 \text{ m/s}}{1.0 \times 10^{-9} \text{ m}}\right) = 2.0 \times 10^{-16} \text{ J}
$$

Assess: This is a very small amount of energy, but it is larger than the energy of a photon in the visible wavelength region.

25.12. Solve: Your mass is, say, $m \approx 70$ kg and your velocity is 1 m/s. Thus, your momentum is $p = mv \approx (70$ $kg(1 \text{ m/s}) = 70 \text{ kg m/s}$. Your de Broglie wavelength is

$$
\lambda = \frac{h}{p} = \frac{6.63 \times 10^{-34} \text{Js}}{70 \text{ kg m/s}} \approx 9 \times 10^{-36} \text{ m}
$$

25.13. Solve: (a) The baseball's momentum is $p = mv = (0.200 \text{ kg})(30 \text{ m/s}) = 6.0 \text{ kg m/s}$. The baseball's de Broglie wavelength is

$$
\lambda = \frac{h}{p} = \frac{6.63 \times 10^{-34} \text{Js}}{6.0 \text{ kg m/s}} = 1.1 \times 10^{-34} \text{ m}
$$

(b) Using $\lambda = h/p = h/mv$, we have

$$
v = \frac{h}{m\lambda} = \frac{6.63 \times 10^{-34} \text{Js}}{(0.200 \text{ kg})(0.20 \times 10^{-9} \text{ m})} = 1.7 \times 10^{-23} \text{ m/s}
$$

25.14. Visualize: We'll employ Equations 25.8 ($\lambda = h/p$) and 25.9 ($E = p^2/2m$) to express the wavelength in terms of kinetic energy.

Solve: First solve Equation 25.9 for $p: p = \sqrt{2mE}$.

$$
\lambda = \frac{h}{p} = \frac{h}{\sqrt{2mE}} = \frac{6.63 \times 10^{-34} \text{ J} \cdot \text{s}}{\sqrt{2(9.11 \times 10^{-31} \text{ kg})(2.4 \times 10^{-19} \text{ J})}} = 1.0 \text{ nm}
$$

Assess: The energy given is about 1.5 eV, which is a reasonable amount of energy. The resulting wavelength is a few to a few dozen times the size of an atom.

25.15. Solve: **(a)** For an electron, the momentum $p = mv = (9.11 \times 10^{-31} \text{ kg})v$. The de Broglie wavelength is

$$
\lambda = \frac{h}{p} = 0.20 \times 10^{-9} \text{ m} \implies 0.20 \times 10^{-9} \text{ m} = \frac{6.63 \times 10^{-34} \text{ Js}}{(9.11 \times 10^{-31} \text{ kg})v} \implies v = 3.6 \times 10^{6} \text{ m/s}
$$

(b) For a proton, $p = mv = (1.67 \times 10^{-9} \text{ kg})v$. The de Broglie wavelength is

$$
\lambda = \frac{h}{p} = 0.20 \times 10^{-9} \text{ m} \Rightarrow 0.20 \times 10^{-9} \text{ m} = \frac{6.63 \times 10^{-34} \text{ Js}}{(1.67 \times 10^{-27} \text{ kg})v} \Rightarrow v = 2.0 \times 10^{3} \text{ m/s}
$$

25.16. Model: The momentum of a wave-like particle has discrete values given by $p_n = n(h/2L)$ where $n =$ $1, 2, 3, \ldots$

Solve: Because we want the smallest box and the momentum of the electron can not exceed a given value, *n* must be minimum. Thus,

$$
p_1 = mv = \frac{h}{2L} \Rightarrow L = \frac{h}{2mv} = \frac{6.63 \times 10^{-34} \text{Js}}{2(9.11 \times 10^{-31} \text{ kg})(10 \text{ m/s})} = 0.036 \text{ mm}
$$

25.17. Model: A confined particle can have only discrete values of energy. **Solve:** From Equation 24.14, the energy of a confined electron is

$$
E_n = \frac{h^2}{8mL^2}n^2 \qquad n = 1, 2, 3, 4, ...
$$

The minimum energy is

$$
E_1 = \frac{h^2}{8mL^2} \Rightarrow L = \frac{h}{\sqrt{8mE_1}} = \frac{6.63 \times 10^{-34} \text{ Js}}{\sqrt{8(9.11 \times 10^{-31} \text{ kg})(1.5 \times 10^{-18} \text{ J})}} = 2.0 \times 10^{-10} \text{ m} = 0.20 \text{ nm}
$$

25.18. Model: Model the 5.0-fm-diameter nucleus as a box of length $L = 5.0$ fm. **Solve:** The proton's energy is restricted to the discrete values

$$
E_n = \frac{h^2}{8mL^2}n^2 = \frac{\left(6.63 \times 10^{-34} \text{ Js}\right)^2 n^2}{8\left(1.67 \times 10^{-27} \text{ kg}\right)\left(5.0 \times 10^{-15} \text{ m}\right)^2} = \left(1.316 \times 10^{-12} \text{ J}\right)n^2
$$

where $n = 1, 2, 3, ...$ For $n = 1, E_1 = 1.3 \times 10^{-12}$ J, for $n = 2, E_2 = (1.316 \times 10^{-12} \text{ J})^2 = 5.3 \times 10^{-12}$ J, and for $n = 3,$ $E_3 = 9E_1 = 1.2 \times 10^{-11}$ J.

25.19. Model: The generalized formula of Balmer predicts a series of spectral lines in the hydrogen spectrum.

Solve: (a) The generalized formula of Balmer

$$
\lambda = \frac{91.18 \text{ m}}{\left(\frac{1}{m^2} - \frac{1}{n^2}\right)}
$$

with $m = 1$ and $n > 1$ accounts for a series of spectral lines. This series is called the Lyman series and the first two members are

$$
\lambda_1 = \frac{91.18 \text{ m}}{\left(1 - \frac{1}{2^2}\right)} = 121.6 \text{ nm} \qquad \lambda_2 = \frac{91.18 \text{ nm}}{\left(1 - \frac{1}{3^2}\right)} = 102.6 \text{ nm}
$$

For $n = 4$ and $n = 5$, $\lambda_3 = 97.3$ nm and $\lambda_4 = 95.0$ nm.

(b) The Lyman series converges when $n \to \infty$. This means $1/n^2 \to 0$ and $\lambda \to 91.18$ nm.

(c) For a diffraction grating, the condition for bright (constructive interference) fringes is $d \sin \theta_p = p\lambda$, where $p =$ 1, 2, 3, ... For first-order diffraction, this equation simplifies to $d \sin \theta = \lambda$. For the first and second members of the Lyman series, the above condition is $d \sin \theta_1 = \lambda_1 = 121.6$ nm and $d \sin \theta_2 = \lambda_2 = 102.6$ nm. Dividing these two equations yields

$$
\sin \theta_2 = \left(\frac{102.6 \text{ nm}}{121.6 \text{ nm}}\right) \sin \theta_1 = (0.84375) \sin \theta_1
$$

The distance from the center to the first maximum is $y = L \tan \theta$. Thus,

$$
\tan \theta_1 = \frac{y_1}{L} = \frac{0.376 \text{ m}}{1.5 \text{ m}} \Rightarrow \theta_1 = 14.072^\circ \Rightarrow \sin \theta_2 = (0.84375) \sin (14.072^\circ) \Rightarrow \theta_2 = 11.84^\circ
$$

Applying the position formula once again,

$$
y_2 = L \tan \theta_2 = (1.5 \text{ m}) \tan (11.84^\circ) = 0.314 \text{ m} = 31.4 \text{ cm}
$$

25.20. Model: The generalized formula of Balmer predicts a series of spectral lines in the hydrogen spectrum.

Solve: (a) The generalized formula of Balmer

$$
\lambda = \frac{91.18 \text{ m}}{\left(\frac{1}{m^2} - \frac{1}{n^2}\right)}
$$

with $m = 3$, and $n > 3$ accounts for a series of spectral lines. This series is called the Paschen series and the wavelengths are

$$
\lambda = \frac{91.18 \text{ nm}}{\left(\frac{1}{3^2} - \frac{1}{n^2}\right)} = \frac{820.62 \text{ n}^2}{n^2 - 9}
$$

The first four members are $\lambda_1 = 1876$ nm, $\lambda_2 = 1282$ nm, $\lambda_3 = 1094$ nm, and $\lambda_4 = 1005$ nm **(b)** The Paschen series converges when $n \rightarrow \infty$. This means

$$
\frac{1}{n^2} \rightarrow 0 \Rightarrow \lambda \rightarrow \frac{91.18 \text{ nm}}{\left(\frac{1}{3}\right)^2} = 820.6 \text{ nm}
$$

(c) For a diffraction grating, the condition for bright (constructive interference) fringes is $d \sin \theta_p = p\lambda$, where $p =$ 1, 2, 3, ... For first-order diffraction, this equation simplifies to $d \sin \theta = \lambda$. For the first and second members of the Paschen series, the condition is $d \sin \theta_1 = \lambda_1$ and $d \sin \theta_2 = \lambda_2$. Dividing these two equations yields

$$
\sin \theta_2 = \sin \theta_1 \left(\frac{\lambda_2}{\lambda_1}\right) = \sin \theta_1 \left(\frac{1282 \text{ nm}}{1876 \text{ nm}}\right) = (0.6834) \sin \theta_1
$$

The distance from the center to the first maximum is $y = L \tan \theta$. Thus,

$$
\tan \theta_1 = \frac{y_1}{L} = \frac{0.607 \text{ m}}{1.5 \text{ m}} = 0.4047 \implies \theta_1 = 22.03^\circ \implies \sin \theta_2 = (0.6834) \sin 22.03^\circ \implies \theta_2 = 14.85^\circ
$$

Applying the position formula once again, $y_2 = L \tan \theta_2 = (1.5 \text{ m}) \tan 14.85^\circ = 0.398 \text{ m} = 39.8 \text{ cm}$

25.21. Model: Use the photon model of light.

Solve: (a) The wavelength is calculated as follows:

$$
E_{\text{gamma}} = hf = h \left(\frac{c}{\lambda} \right) \Longrightarrow \lambda = \frac{\left(6.63 \times 10^{-34} \text{ Js} \right) \left(3.0 \times 10^8 \text{ m/s} \right)}{1.0 \times 10^{-13} \text{ J}} = 2.0 \times 10^{-12} \text{ m}
$$

(b) The energy of a visible-light photon of wavelength 500 nm is

$$
E_{\text{visible}} = h \left(\frac{c}{\lambda} \right) = \frac{\left(6.63 \times 10^{-34} \text{ Js} \right) \left(3.0 \times 10^8 \text{ m/s} \right)}{500 \times 10^{-9} \text{ m}} = 3.978 \times 10^{-19} \text{ J}
$$

The number of photons *n* such that $E_{\text{gamma}} = nE_{\text{visible}}$ is

$$
n = \frac{E_{\text{gamma}}}{E_{\text{visible}}} = \frac{1.0 \times 10^{-13} \text{ J}}{3.978 \times 10^{-19} \text{ J}} = 2.5 \times 10^5
$$

25.22. Model: Use the photon model.

Solve: The energy of a 1000 kHz photon is

$$
E_{\text{photon}} = hf = (6.63 \times 10^{-34} \text{ Js})(1000 \times 10^3 \text{ Hz}) = 6.63 \times 10^{-28} \text{ J}
$$

The energy transmitted each second is 20×10^3 J. The number of photons transmitted each second is 20×10^3 J/6.63 $\times 10^{-28}$ J = 3.0 $\times 10^{31}$.

25.23. Model: Use the photon model for the laser light. **Solve: (a)** The energy is

$$
E_{\text{photon}} = hf = h \left(\frac{c}{\lambda} \right) = \left(6.63 \times 10^{-34} \text{ Js} \right) \left(\frac{3 \times 10^8 \text{ m/s}}{633 \times 10^{-9} \text{ m}} \right) = 3.1 \times 10^{-19} \text{ J}
$$

(b) The energy emitted each second is 1.0×10^{-3} J. The number of photons emitted each second is 1.0×10^{-3} J/3.14 $\times 10^{-19}$ J = 3.2 $\times 10^{15}$.

25.24. Model: Use the photon model for the incandescent light.

Solve: The photons travel in all directions. At a distance of *r* from the light bulb, the photons spread over a sphere of surface area $4\pi r^2$. The number of photons per second per unit area at the location of your retina is

$$
\frac{3 \times 10^{18} \text{ s}^{-1}}{4 \pi (10 \times 10^3 \text{ m})^2} = 2.387 \times 10^9 \text{ s}^{-1} \text{ m}^{-2}
$$

The number of photons that enter your pupil per second is

$$
2.387 \times 10^9 \text{ s}^{-1} \text{m}^{-2} \times \pi \left(3.5 \times 10^{-3} \text{ m}\right)^2 = 9.2 \times 10^4 \text{ s}^{-1}
$$

25.25. Model: Use the photon model of light and the Bragg condition for diffraction.

Solve: The Bragg condition for the reflection of x-rays from a crystal is $2d\cos\theta_m = m\lambda$. To determine the angles of incidence θ_m , we need to first calculate λ . The wavelength is related to the photon's energy as $E = hc/\lambda$. Thus,

$$
\lambda = \frac{hc}{E} = \frac{(6.63 \times 10^{-34} \text{ Js})(3.0 \times 10^8 \text{ m/s})}{1.50 \times 10^{-15} \text{ J}} = 1.326 \times 10^{-10} \text{ m}
$$

From the Bragg condition,

$$
\theta_m = \cos^{-1}\left(\frac{m\lambda}{2d}\right) = \cos^{-1}\left[\frac{(1.326 \times 10^{-10} \text{ m})m}{2(0.21 \times 10^{-9} \text{ m})}\right] = \cos^{-1}(0.3157m) \Rightarrow \theta_1 = \cos^{-1}(0.3157) = 71.6^{\circ}
$$

Likewise, $\theta_2 = \cos^{-1}(0.3157 \times 2) = 50.8^\circ$ and $\theta_3 = 18.7^\circ$. Note that $\theta_4 = \cos^{-1}(0.3157 \times 4)$ is not allowed because the cos θ cannot be larger than 1. Thus, the x-rays will be diffracted at angles of incidence equal to 18.7°, 50.8°, and 71.6°.

25.26. Model: The angles for which diffraction from parallel planes occurs satisfy the Bragg condition. **Solve:** We cannot assume that these are the first and second order diffractions. The Bragg condition is $2d\cos\theta_m = m\lambda$. We have

$$
2d\cos 45.6^\circ = m\lambda \qquad 2d\cos 21.0^\circ = (m+1)\lambda
$$

Notice that θ*m* decreases as *m* increases, so 21.6° corresponds to the larger value of *m*. Dividing these two equations,

$$
\frac{\cos 45.6^{\circ}}{\cos 21.0^{\circ}} = \frac{m}{m+1} = 0.7494 \Rightarrow m = 3
$$

Thus these are the third and fourth order diffractions. Substituting into the Bragg condition,

$$
d = \frac{3 \times 0.0700 \times 10^{-9} \text{ m}}{2 \cos 45.6^{\circ}} = 1.50 \times 10^{-10} \text{ m} = 0.150 \text{ nm}
$$

25.27. Model: The x-ray diffraction angles satisfy the Bragg condition.

Solve: (a) The Bragg condition ($2d \cos \theta_m = m\lambda$) for normal incidence, $\theta_m = 0^\circ$, simplifies to $2d = m\lambda$. For a thin film of a material on a substrate where $n_{air} < n_{material} < n_{substrate}$, constructive interference between the two reflected waves occurs when $2d = m\lambda$, where λ is the wavelength inside the material.

(b) From a thin film with a period of 1.2 nm, that is, with $d = 1.2$ nm, the two longest x-ray wavelengths that will reflect at normal incidence are

$$
\lambda_1 = \frac{2d}{1} \qquad \lambda_2 = \frac{2d}{2}
$$

This means that $\lambda_1 = 2(1.2 \text{ nm}) = 2.4 \text{ nm}$ and $\lambda_2 = 1.2 \text{ nm}$.

25.28. Solve: A small fraction of the light wave of an appropriate wavelength is reflected from each little "bump" in the refractive index. These little bumps act like the atomic planes in a crystal. The light will be strongly reflected (and hence blocked in transmission) if it satisfies the Bragg condition at normal incidence $(\theta = 0)$.

$$
2d = m\lambda_{\text{glass}} = \frac{m\lambda}{n_{\text{glass}}} \Rightarrow \lambda = \frac{2dn_{\text{glass}}}{m} = \frac{2(0.45 \times 10^{-6} \text{ m})(1.50)}{1} = 1.35 \text{ }\mu\text{m}
$$

25.29. Model: Particles have a de Broglie wavelength given by $\lambda = h/p$. The wave nature of the particles causes an interference pattern in a double-slit apparatus.

Solve: (a) Since the speed of the neutron and electron are the same, the neutron's momentum is

$$
p_{n} = m_{n}v_{n} = \frac{m_{n}}{m_{e}}m_{e}v_{n} = \frac{m_{n}}{m_{e}}m_{e}v_{e} = \frac{m_{n}}{m_{e}}p_{e}
$$

where m_n and m_e are the neutron's and electron's masses. The de Broglie wavelength for the neutron is

$$
\lambda_{\rm n} = \frac{h}{p_{\rm n}} = \frac{h}{p_{\rm e}} \frac{p_{\rm e}}{p_{\rm n}} = \lambda_{\rm e} \frac{m_{\rm e}}{m_{\rm n}}
$$

From Section 22.2 on double-slit interference, the fringe spacing is $\Delta y = \lambda L / d$. Thus, the fringe spacing for the electron and neutron are related by

$$
\Delta y_{\rm n} = \frac{\lambda_{\rm n}}{\lambda_{\rm e}} \Delta y_{\rm e} = \frac{m_{\rm e}}{m_{\rm n}} \Delta y_{\rm e} = \left(\frac{9.11 \times 10^{-31} \text{ kg}}{1.67 \times 10^{-27} \text{ kg}}\right) (1.5 \times 10^{-3} \text{ m}) = 8.18 \times 10^{-7} \text{ m} = 0.818 \ \mu\text{m}
$$

(b) If the fringe spacing has to be the same for the neutrons and the electrons,

$$
\Delta y_e = \Delta y_n \implies \lambda_e = \lambda_n \implies \frac{h}{m_e v_e} = \frac{h}{m_n v_n} \implies v_n = v_e \frac{m_e}{m_n} = \left(2.0 \times 10^6 \text{ m/s}\right) \left(\frac{9.11 \times 10^{-31} \text{ kg}}{1.67 \times 10^{-27} \text{ kg}}\right) = 1.1 \times 10^3 \text{ m/s}
$$

25.30. Model: Electrons have a de Broglie wavelength given by $\lambda = h/p$. The wave nature of the electrons causes a diffraction pattern.

Solve: The width of the central maximum of a single-slit diffraction pattern is given by Equation 22.22:

$$
w = \frac{2\lambda L}{a} = \frac{2Lh}{ap} = \frac{2Lh}{amv} = \frac{2(1.0 \text{ m})(6.63 \times 10^{-34} \text{ Js})}{(1.0 \times 10^{-6} \text{ m})(9.11 \times 10^{-31} \text{ kg})(1.5 \times 10^{6} \text{ m/s})} = 9.7 \times 10^{-4} \text{ m} = 0.97 \text{ mm}
$$

25.31. Model: Neutrons have a de Broglie wavelength given by $\lambda = h/p$. The wave nature of the neutrons causes a double-slit interference pattern.

Solve: Measurements show that the spacing between the $m = 1$ and $m = -1$ peaks is 1.4 times as long as the length of the reference bar, which gives the real fringe separation Δ*y* = 70 μm. Similarly, the spacing between the $m = 2$ and $m = -2$ is 2.8 times as long as the length of the reference bar and yields $\Delta y = 70 \mu m$.

The fringe separation in a double-slit experiment is $\Delta y = \lambda L/d$. Hence,

$$
\lambda = \frac{\Delta y}{L} \Rightarrow \frac{h}{p} = \frac{h}{mv} = \frac{\Delta y}{L} \Rightarrow v = \frac{hL}{\Delta y \, md} = \frac{(6.63 \times 10^{-34} \, \text{Js})(3.0 \, \text{m})}{(70 \times 10^{-6} \, \text{m})(1.67 \times 10^{-27} \, \text{kg})(0.10 \times 10^{-3} \, \text{m})} = 170 \, \text{m/s}
$$

25.32. Model: Electrons have a de Broglie wavelength given by $\lambda = h/p$.

Visualize: Please refer to Figure 25.11. Notice that a scattering angle $\phi = 60^\circ$ corresponds to an angle of incidence $\theta = 30^\circ$.

Solve: Equation 25.6 describes the Davisson-Germer experiment: $D\sin(2\theta_m) = m\lambda$. Assuming $m = 1$, this equation simplifies to $D \sin 2\theta = \lambda$. Using $\lambda = h/mv$, we have

 $(9.11\times10^{-31} \text{ kg})$ $(4.30\times10^{6} \text{ m/s})$ sin (60°) $D = \frac{h}{m v \sin 2\theta} = \frac{6.63 \times 10^{-34} \text{ Js}}{(9.11 \times 10^{-31} \text{ kg})(4.30 \times 10^6 \text{ m/s}) \sin(60^\circ)} = 1.95 \times 10^{-10} \text{ m} = 0.195 \text{ nm}$ **25.33. Model:** A confined particle can have only discrete values of energy. **Solve: (a)** Equation 25.14 simplifies to

$$
E_n = \frac{h^2}{8mL^2} n^2 = \frac{\left(6.63 \times 10^{-34} \text{ Js}\right)^2}{8\left(9.11 \times 10^{-31} \text{ kg}\right)\left(0.70 \times 10^{-9} \text{ m}\right)^2} = \left(1.231 \times 10^{-19} \text{ J}\right) n^2
$$

Thus, $E_1 = (1.231 \times 10^{-19} \text{ J})(1^2) = 1.2 \times 10^{-19} \text{ J}, E_2 = (1.231 \times 10^{-19} \text{ J})(2^2) = 4.9 \times 10^{-19} \text{ J}, \text{ and } E_3 = 1.1 \times 10^{-18} \text{ J}.$ **(b)** The energy is $E_2 - E_1 = 4.9 \times 10^{-19}$ J -1.2×10^{-19} J = 3.7×10^{-19} J.

(c) Because energy is conserved, the photon will carry an energy of $E_2 - E_1 = 3.69 \times 10^{-19}$ J. That is,

$$
E_2 - E_1 = E_{\text{photon}} = hf = \frac{hc}{\lambda} \Longrightarrow \lambda = \frac{hc}{E_2 - E_1} = \frac{(6.63 \times 10^{-34} \text{ Js})(3.0 \times 10^8 \text{ m/s})}{3.69 \times 10^{-19} \text{ J}} = 540 \text{ nm}
$$

25.34. Model: A particle confined in a one-dimensional box has discrete energy levels. **Solve:** (a) Equation 24.14 for the $n = 1$ state is

$$
E_n = \frac{h^2}{8mL^2} = \frac{\left(6.63 \times 10^{-34} \text{ Js}\right)^2}{8\left(10 \times 10^{-3} \text{ kg}\right)\left(0.10 \text{ m}\right)^2} = 5.5 \times 10^{-64} \text{ J}
$$

The minimum energy of the Ping-Pong ball is $E_1 = 5.5 \times 10^{-64}$ J.

(b) The speed is calculated as follows:

$$
E_1 = 5.50 \times 10^{-64} \text{ J} = \frac{1}{2} m v^2 = \frac{1}{2} \left(10 \times 10^{-3} \text{ kg} \right) v^2 \implies v = \sqrt{\frac{2 \left(5.50 \times 10^{-64} \text{ J} \right)}{10 \times 10^{-3} \text{ kg}}} = 3.3 \times 10^{-31} \text{ m/s}
$$

25.35. Model: A particle confined in a one-dimensional box has discrete energy levels. **Solve:** Using Equation 24.14 for $n = 1$ and 2,

$$
E_2 - E_1 = \frac{h^2}{8mL^2} \left(2^2 - 1^2 \right) \Rightarrow 1.0 \times 10^{-19} \text{ J} = \frac{\left(6.63 \times 10^{-34} \text{ Js} \right)^2}{8 \left(9.11 \times 10^{-31} \text{ kg} \right) L^2} (3) = \frac{1.809 \times 10^{-37} \text{ Jm}^2}{L^2}
$$

$$
\Rightarrow L = \sqrt{\frac{1.809 \times 10^{-37} \text{ Jm}^2}{1.0 \times 10^{-19} \text{ J}}} = 1.3 \times 10^{-9} \text{ m} = 1.3 \text{ nm}
$$

25.36. Visualize: From the figure we see that the wavelength is 2.0 nm. We'll employ Equations 25.8 $(\lambda = h/p)$ and 25.9 $(E = p^2/2m)$ to express the kinetic energy in terms of wavelength. **Solve:**

$$
E = \frac{p^2}{2m} = \frac{(h/\lambda)^2}{2m} = \frac{(6.63 \times 10^{-34} \,\text{J} \cdot \text{s}/2.0 \text{ nm})^2}{2(9.11 \times 10^{-31} \,\text{kg})} = 6.0 \times 10^{-20} \,\text{J}
$$

Assess: This energy is a little less than one eV, which is reasonable.

25.37. Visualize: The strategy is to take ratios to find *n* and then plug it back in to find *L*. **Solve:**

$$
\frac{E_{n+1}}{E_n} = \frac{\frac{h^2}{8mL^2}(n+1)^2}{\frac{h^2}{8mL^2}(n)^2} = \frac{(n+1)^2}{(n)^2} = \frac{6.4 \times 10^{-13} \text{ J}}{3.6 \times 10^{-13} \text{ J}}
$$

Cancel $\times 10^{-13}$ J and take square roots.

$$
\frac{n+1}{n} = \sqrt{\frac{6.4}{3.6}} = \frac{4}{3} \qquad \Rightarrow \qquad n = 3
$$

So $E_3 = 3.6 \times 10^{-13}$ J. Now solve for *L*.

$$
L = \sqrt{\frac{h^2 n^2}{8mE_n}} = \frac{hn}{\sqrt{8mE_n}} = \frac{(6.63 \times 10^{-34} \text{ J} \cdot \text{s})(3)}{\sqrt{8(1.67 \times 10^{-27} \text{ kg})(3.6 \times 10^{-13} \text{ J})}} = 29 \text{ fm}
$$

Assess: This is not an atomic-sized box, but a nuclear-sized box; that's OK for neutrons.

25.38. Model: The allowed energies of a particle of mass *m* in a two-dimensional square box of side *L* are

$$
E_{nm}=\frac{h^2}{8mL^2}\left(n^2+m^2\right)
$$

Solve: (a) The minimum energy for a particle is for $n = m = 1$:

$$
E_{\min} = E_{11} = \frac{h^2}{8mL^2} \left(1^2 + 1^2 \right) = \frac{h^2}{4mL^2}
$$

(b) The five lowest allowed energies are E_{min} , $\frac{5}{2}E_{\text{min}}$ (for $n = 1$, $m = 2$ and $n = 2$, $m = 1$), $4E_{\text{min}}$ (for $n = 2$, $m = 2$), 5*E*_{min} (for $n = 1$, $m = 3$ and $n = 3$, $m = 1$), and $\frac{13}{2}$ *E*_{min} (for $n = 2$, $m = 3$ and $n = 3$, $m = 2$).

25.39. Model: A particle confined in a one-dimensional box of length *L* has the discrete energy levels given by Equation 24.14.

Solve: (a) Since the energy is entirely kinetic energy,

$$
E_n = \frac{h^2}{8mL^2}n^2 = \frac{p^2}{2m} = \frac{1}{2}mv_n^2 \Rightarrow v_n = \frac{h}{2mL}n \qquad n = 1, 2, 3, ...
$$

(b) The first allowed velocity is

$$
v_1 = \frac{6.63 \times 10^{-34} \text{Js}}{2(9.11 \times 10^{-31} \text{ kg})(0.20 \times 10^{-9} \text{ m})} = 1.82 \times 10^6 \text{ m/s}
$$

For $n = 2$ and $n = 3$, $v_2 = 3.64 \times 10^6$ m/s and $v_3 = 5.46 \times 10^6$ m/s.

25.40. Model: Sets of parallel planes in a crystal diffract x-rays.

Visualize: Please refer to Figure CP25.40.

Solve: The Bragg diffraction condition is $2d\cos\theta_m = m\lambda$, where *d* is the interplanar separation. Because smaller *m* values correspond to higher angles of incidence, the diffraction angle of 71.3° in the x-ray intensity plot must correspond to $m = 1$. This means

$$
2d\cos 71.3^\circ = 1(0.10 \times 10^{-9} \text{ m}) \Rightarrow d = \frac{0.10 \times 10^{-9} \text{ m}}{2(\cos 71.3^\circ)} = 1.56 \times 10^{-10} \text{ m}
$$

The cosines of the three angles in the x-ray intensity plot are $\cos 71.3^\circ = 0.321$, $\cos 50.1^\circ = 0.642$, and $\cos 15.8^\circ = 0.962$. These are in the ratio 1 : 2 : 3, which tells us that these are the $m = 1, 2$, and 3 diffraction peaks from a single set of planes with $d = 0.156$ nm.

We can see from the figure that the atomic spacing *D* of this crystal is related to the interplanar separation *d* by

$$
D = \frac{d}{\sin 60^{\circ}} = \frac{0.156 \text{ nm}}{\sin 60^{\circ}} = 0.18 \text{ nm}
$$

25.41. Model: Sets of parallel planes in a crystal diffract x-rays.

Visualize: Please refer to Figure 25.7.

Solve: (a) The Bragg diffraction condition is $2d\cos\theta_m = m\lambda$. The plane spacing is $d_A = 0.20$ nm and the x-ray wavelength is $\lambda = 0.12$ nm. Thus

$$
\cos \theta_m = \frac{m\lambda}{2d_A} = \frac{m(0.12 \times 10^{-9} \text{ m})}{2(0.20 \times 10^{-9} \text{ m})} = (0.3) m \Rightarrow \theta_{A1} = \cos^{-1}(0.3) = 72.5^{\circ}
$$

Likewise for $m = 2$ and $m = 3$, $\theta_{A2} = \cos^{-1}(0.6) = 53.1^\circ$ and $\theta_{A3} = \cos^{-1}(0.9) = 25.8^\circ$. These three angles for the x-ray diffraction peaks match the peaks shown in Figure 25.7c.

(b) The new interplaner spacing is $d_B = d_A / \sqrt{2} = 0.141$ nm (see Figure 25.7b). The Bragg condition for the tilted atomic planes becomes

$$
\cos \theta_m = \frac{m\lambda}{2d_\text{B}} = 0.4243m
$$

For $m = 1$, $\theta_{\text{B1}} = \cos^{-1}(0.4243) = 64.9^{\circ}$. For $m = 2$, $\theta_{\text{B2}} = \cos^{-1}(0.8486) = 31.9^{\circ}$.

(c) The crystal is already tipped by 45° to get the tilted planes (see Figure 25.7b). So, for $m = 1$, $\theta_1 = 64.9^{\circ} - 45^{\circ} = 19.9^{\circ}$. $\theta = 64.9^{\circ} + 45^{\circ} = 109.9^{\circ}$ also, but we can't see beyond 90°. For $m = 2$, $\theta_2 = 31.9^\circ + 45^\circ = 76.9^\circ$. These two angles match the angles in the diffraction peaks of the tilted planes.

25.42. Model: This is an integrated problem that uses concepts from Chapter 22. There are two *L*'s in the problem: *L* in Chapter 22 refers to the screen distance from the slits, and the *L* we want here is the length of the box. The wavelength of the neutron determined by the two-slit pattern is the same as the wavelength in the confined box.

Visualize: The figure shows $L_{\text{box}} = 2\lambda$.

We also need Equation 22.6: $y_m = \frac{m\lambda L_{\text{screen}}}{d}$. Also from the figure we see that $y_2 = 0.20 \times 10^{-3}$ m. We are given $L_{\text{screen}} = 2.0 \text{ m} \text{ and } d = 15 \times 10^{-6} \text{ m}.$

Solve: Solve Equation 22.6 for λ .

$$
\lambda = \frac{dy_m}{mL_{\text{screen}}}
$$

$$
L_{\text{box}} = 2\lambda = 2 \frac{dy_m}{mL_{\text{screen}}} = 2 \frac{(15 \times 10^{-6} \text{ m})(0.20 \times 10^{-3} \text{ m})}{(2)(2.0 \text{ m})} = 1.5 \text{ nm}
$$

Assess: The two pieces of this problem fit together and make sense together.

25.43. Model: Electrons have a de Broglie wavelength given by $\lambda = h/p$. Trapped electrons in the confinement layer behave like a de Broglie wave in a closed-closed tube or like a string fixed at both ends.

Solve: (a) The four longest standing-wave wavelengths in the layer are $\lambda = 2L$, L , $\frac{2}{3}L$, and $\frac{1}{2}L$. This follows from the general relation for closed-closed tubes: $\lambda = 2L/n$. Thus, $\lambda = 10.0$ nm, 5.00 nm, 3.33 nm, and 2.50 nm. **(b)** We have

$$
p = mv = \frac{h}{\lambda} \Rightarrow v = \frac{h}{m\lambda} = \frac{6.63 \times 10^{-34} \text{Js}}{(9.11 \times 10^{-31} \text{ kg})\lambda} = \frac{0.7278 \times 10^{-3} \text{ m}^2/\text{s}}{\lambda}
$$

Using the above four longest values of λ we get the four smallest values of ν . Thus,

$$
v_1 = \frac{0.7278 \times 10^{-3} \text{ m}^2/\text{s}}{10.0 \times 10^{-9} \text{ m}} = 7.28 \times 10^4 \text{ m/s}
$$

 $v_2 = 1.46 \times 10^5$ m/s, $v_3 = 2.18 \times 10^5$ m/s, and $v_4 = 2.91 \times 10^5$ m/s.

25.44. Model: As light is diffracted by matter, matter can also be diffracted by light. **Solve:** The de Broglie wavelength of the sodium atoms is

$$
\lambda = \frac{h}{p} = \frac{h}{mv} = \frac{6.63 \times 10^{-34} \text{ Js}}{(3.84 \times 10^{-26} \text{ kg})(50 \text{ m/s})} = 3.45 \times 10^{-10} \text{ m}
$$

The slit spacing of the "diffraction grating" is $d = \frac{1}{2} \lambda_{\text{laser}} = \frac{1}{2} 600 \text{ nm} = 300 \text{ nm}$. Using the diffraction grating equation with $m = 1$, we have

$$
d\sin\theta = (1)\lambda \Rightarrow \sin\theta = \frac{\lambda}{d} = 1.151 \times 10^{-3} \approx 1.2 \times 10^{-3}
$$

In the small-angle approximation, $\sin \theta \approx \tan \theta = y/L$. We get

$$
y = L\sin\theta = (1.0 \text{ m})(1.151 \times 10^{-3}) = 1.2 \text{ mm}
$$